

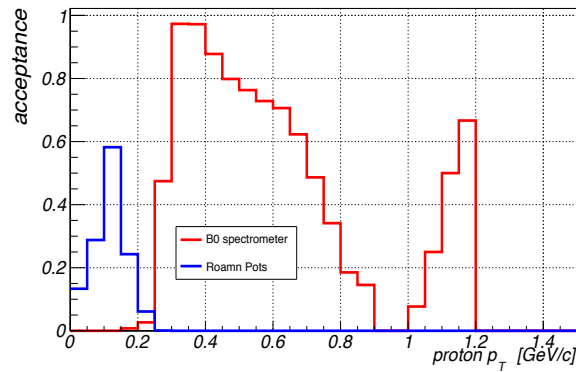
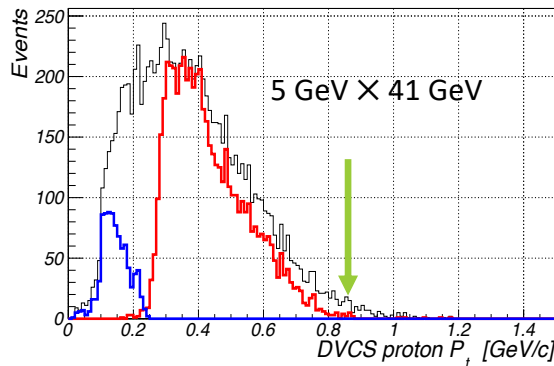
The eRHIC IR and diffractive eA physics

WAN CHANG

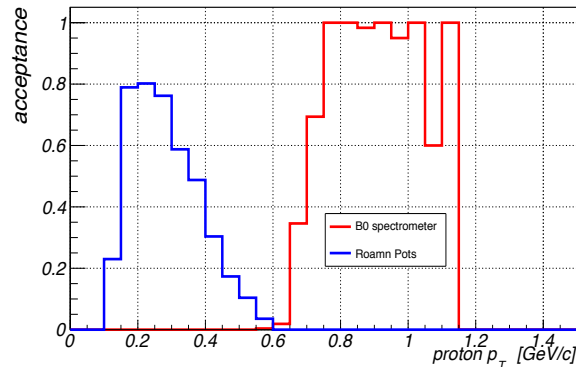
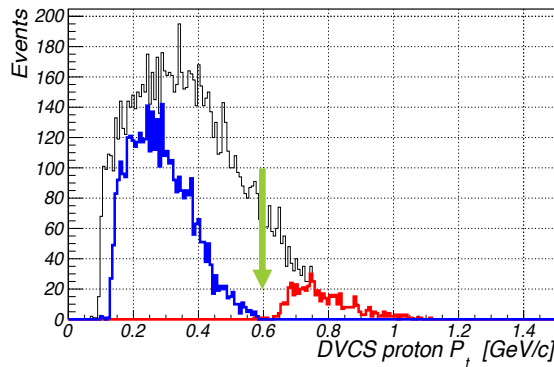
2019/04/03

Note: These plots are for the IR design from the 2018 eRHIC pre-CDR and that we are working on a revised IR design, which will be discussed in the updated pre_CDR

The p_T acceptance of proton from DVCS process

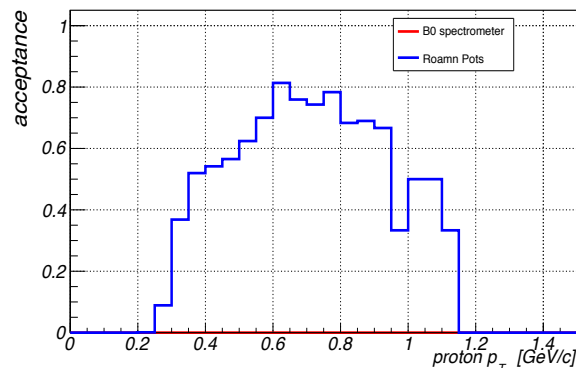
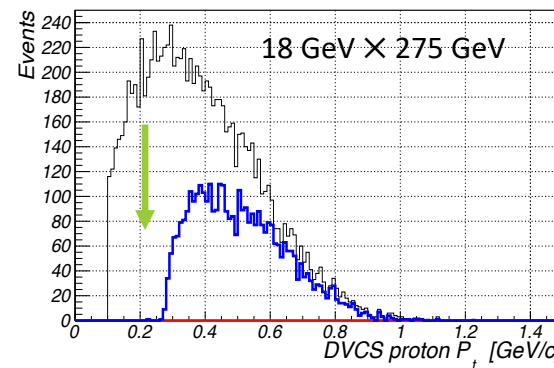


Protons with scattering angle from 0-5mrad are detected by Roman Pots, while the range from 7 to 20mrad is covered by B0 spectrometer



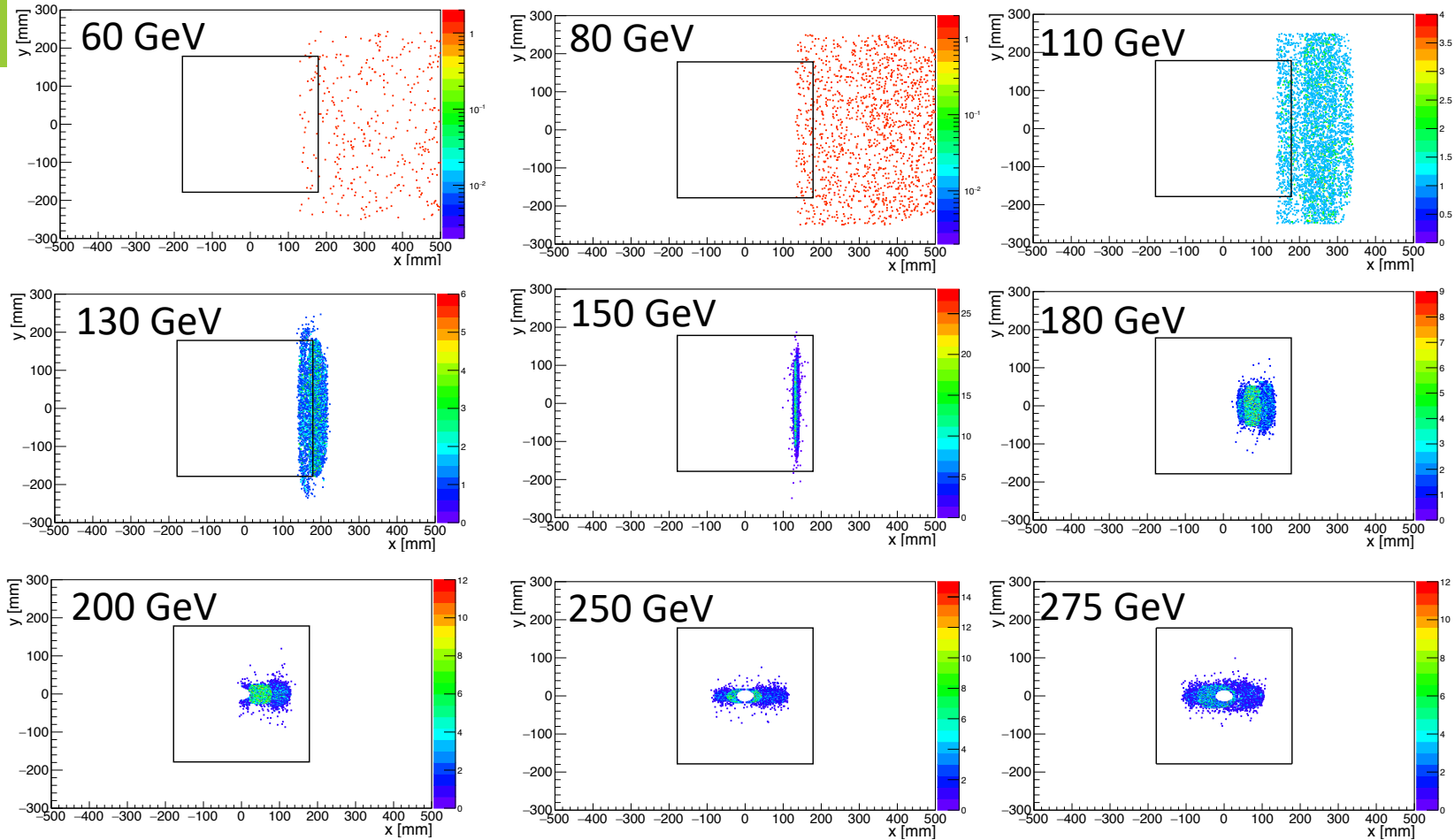
The limiting factors:

- ❖ the inner dimension of the vacuum chamber and magnet apertures
- ❖ beam pipe
- ❖ 10σ away cut from beam line



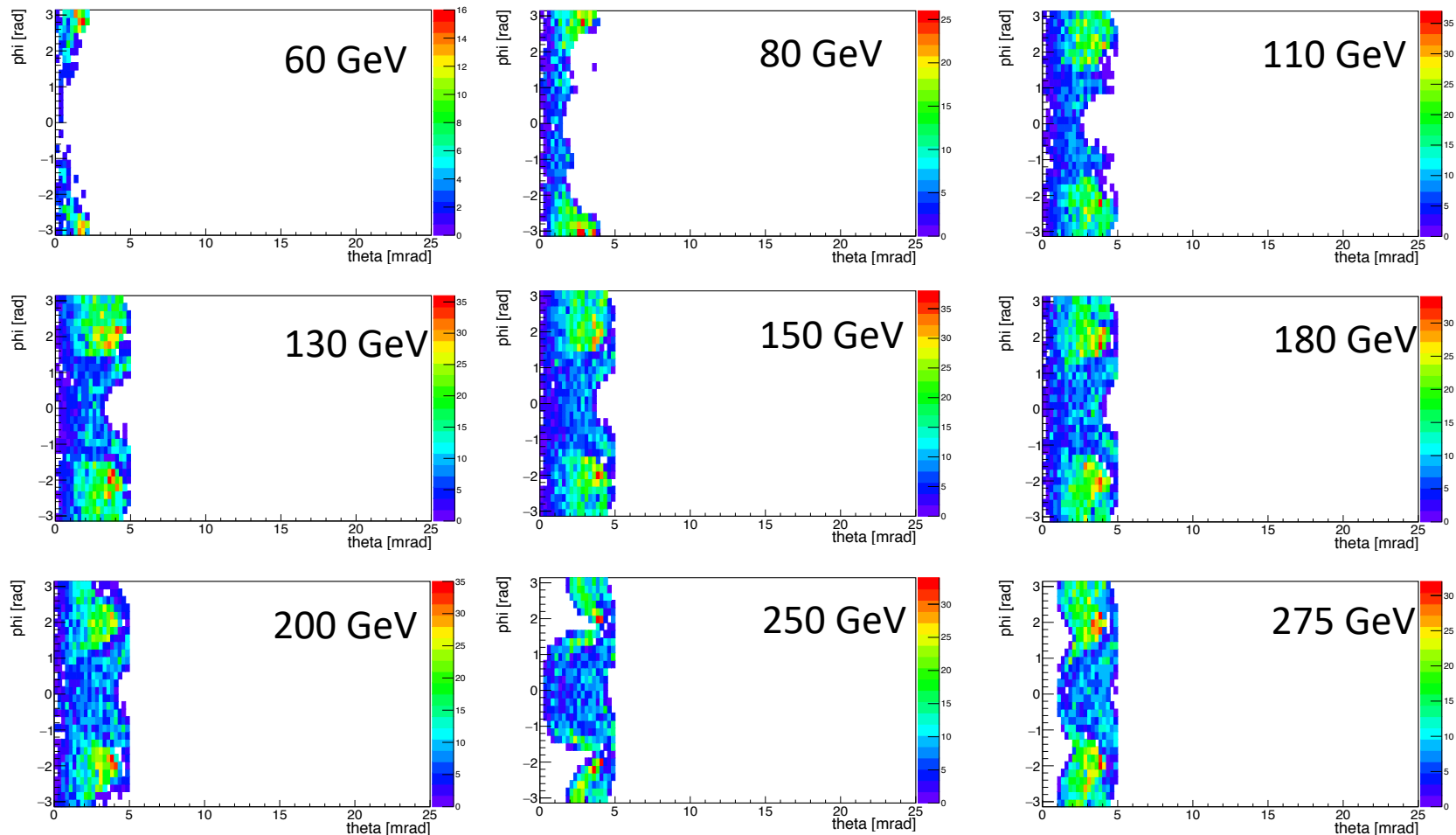
— Roman Pots
— B0 spectrometer

X_Y distribution at Roman Pots position(fix momentum & theta 0_5mrad)



- ❖ Z of Roman Pots: 28.1m; Possible real size: 357mm * 357mm
- ❖ As the p decreases, the rigidity goes down, and the proton bends more
- ❖ The hole in the middle of higher momentum distribution is due to 10σ away cut from the beam line

Theta VS Phi at Roman Pots position(fix momentum & theta 0_5mrad)



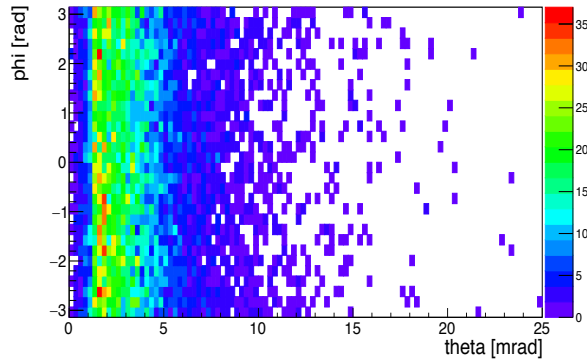
- ❖ The lose particles in low momentum is because the magnets bend them away
- ❖ The gap at $\theta \sim 0$ for high momentum protons is due to the cut of 10σ away from the beam line

ePb Data sample

- ❑ e+Pb(BeAGLE)
- ❑ 18×110 (GeV)
- ❑ $1 < Q^2 < 10$
- ❑ $0.01 < y < 0.95$
- ❑ 10K events

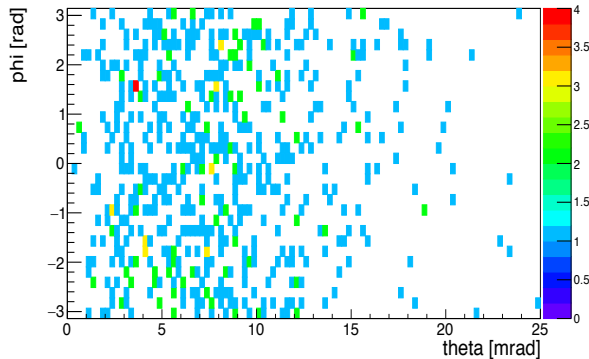
MC generated protons theta VS phi

MC generated particle

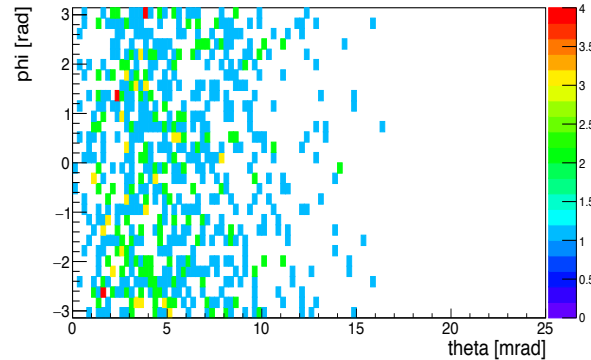


The protons are uniformly distributed in phi

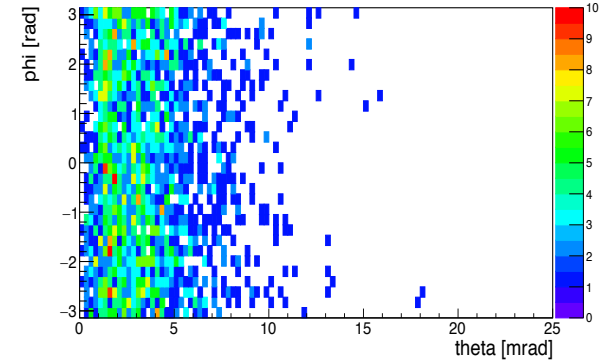
MC generated particle with momentum of 0_60 GeV



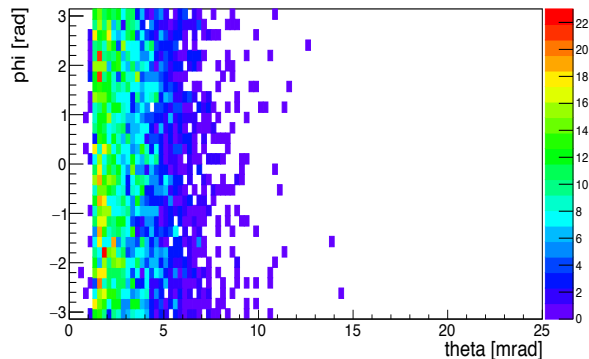
MC generated particle with momentum of 60_80 GeV



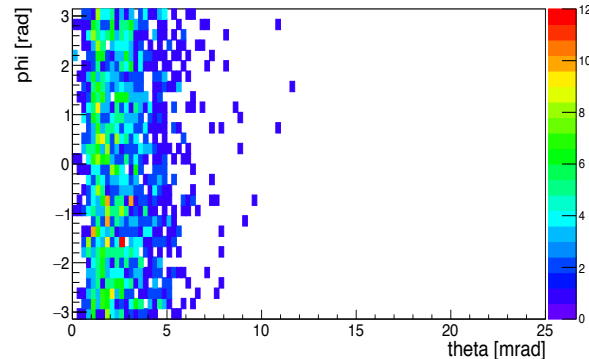
MC generated particle with momentum of 80_100 GeV



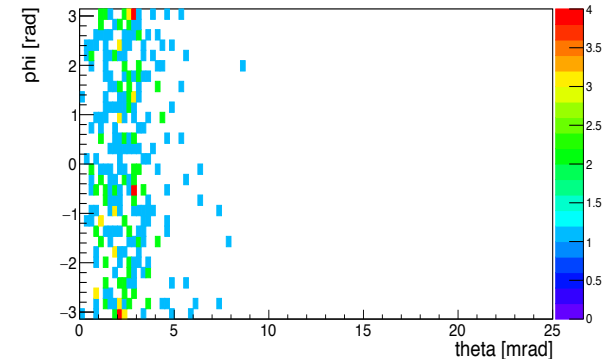
MC generated particle with momentum of 100_120 GeV



MC generated particle with momentum of 120_140 GeV

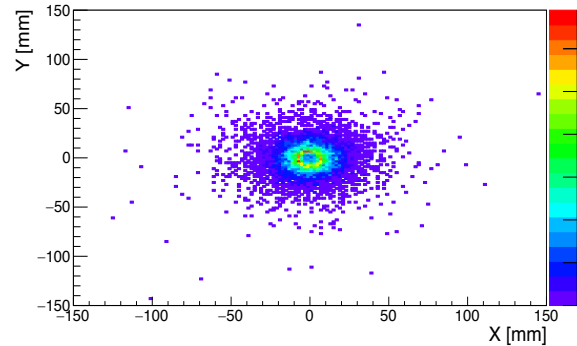


MC generated particle with momentum of 140_200 GeV



MC generate particle X VS Y

X_Y distribution of MC generated



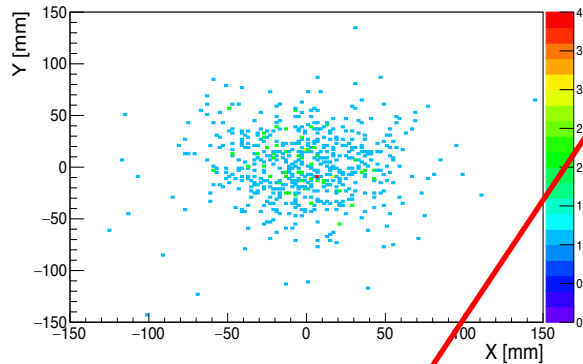
$Z = 4.9\text{m}$ ($b_0: 5.0\text{m}-6.2\text{m}$);

$X = Z * \tan(\theta) * \cos(\phi)$;

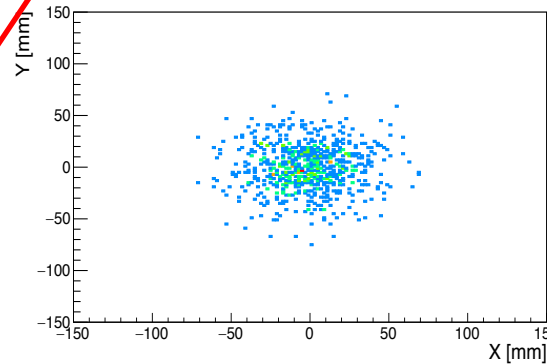
$Y = Z * \tan(\theta) * \sin(\phi)$

Due to the interactions between the outgoing protons and the nuclear remnants, the proton needs some minimum momentum to escape from its mother, so there is some phase space where one doesn't see the protons ($91\text{ GeV} < p_z < 129\text{ GeV}$).

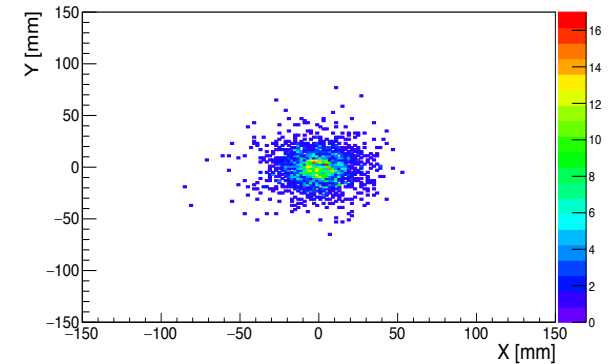
MC generated particle with momentum of 0_60 GeV



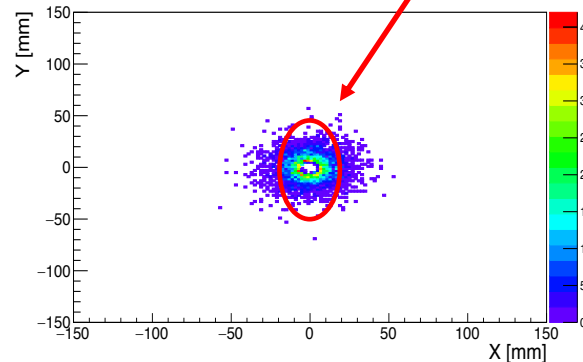
MC generated particle with momentum of 60_80 GeV



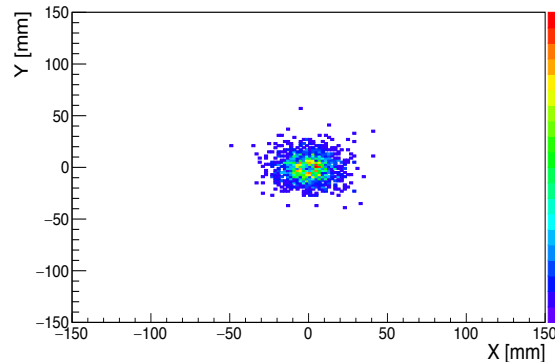
MC generated particle with momentum of 80_100 GeV



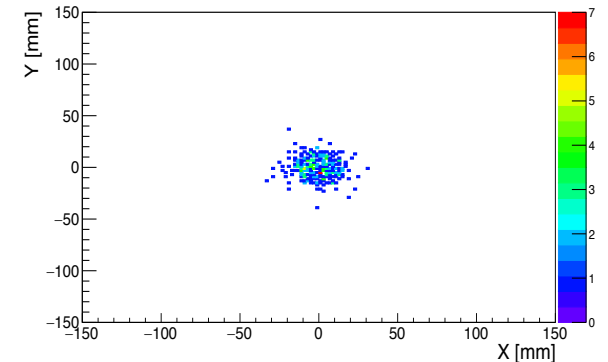
MC generated particle with momentum of 100_120 GeV



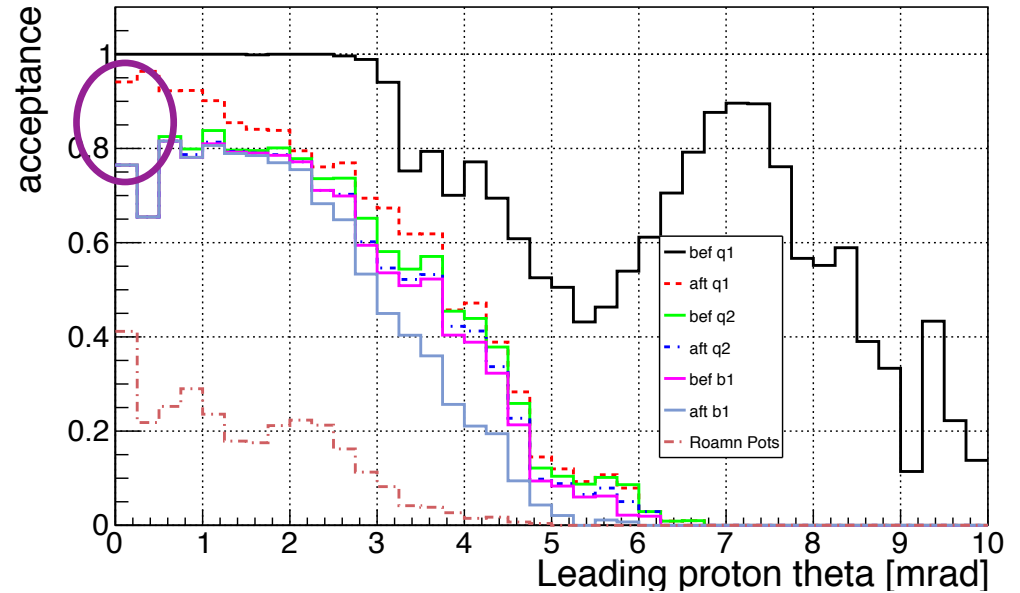
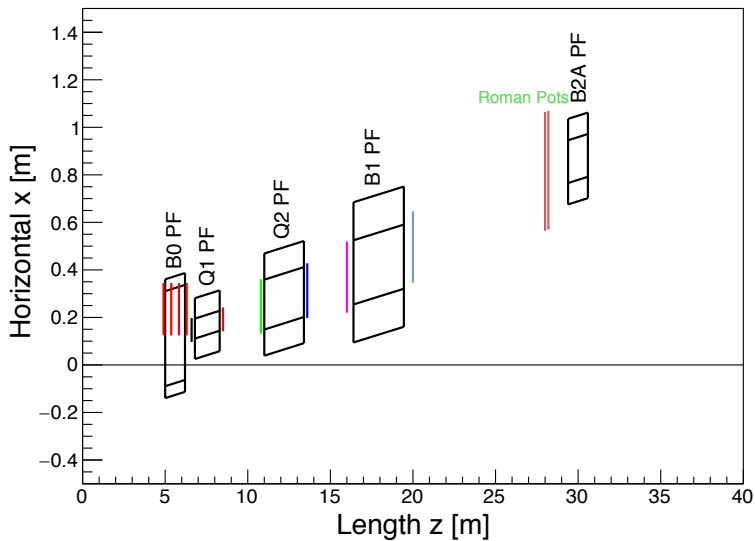
MC generated particle with momentum of 120_140 GeV



MC generated particle with momentum of 140_200 GeV



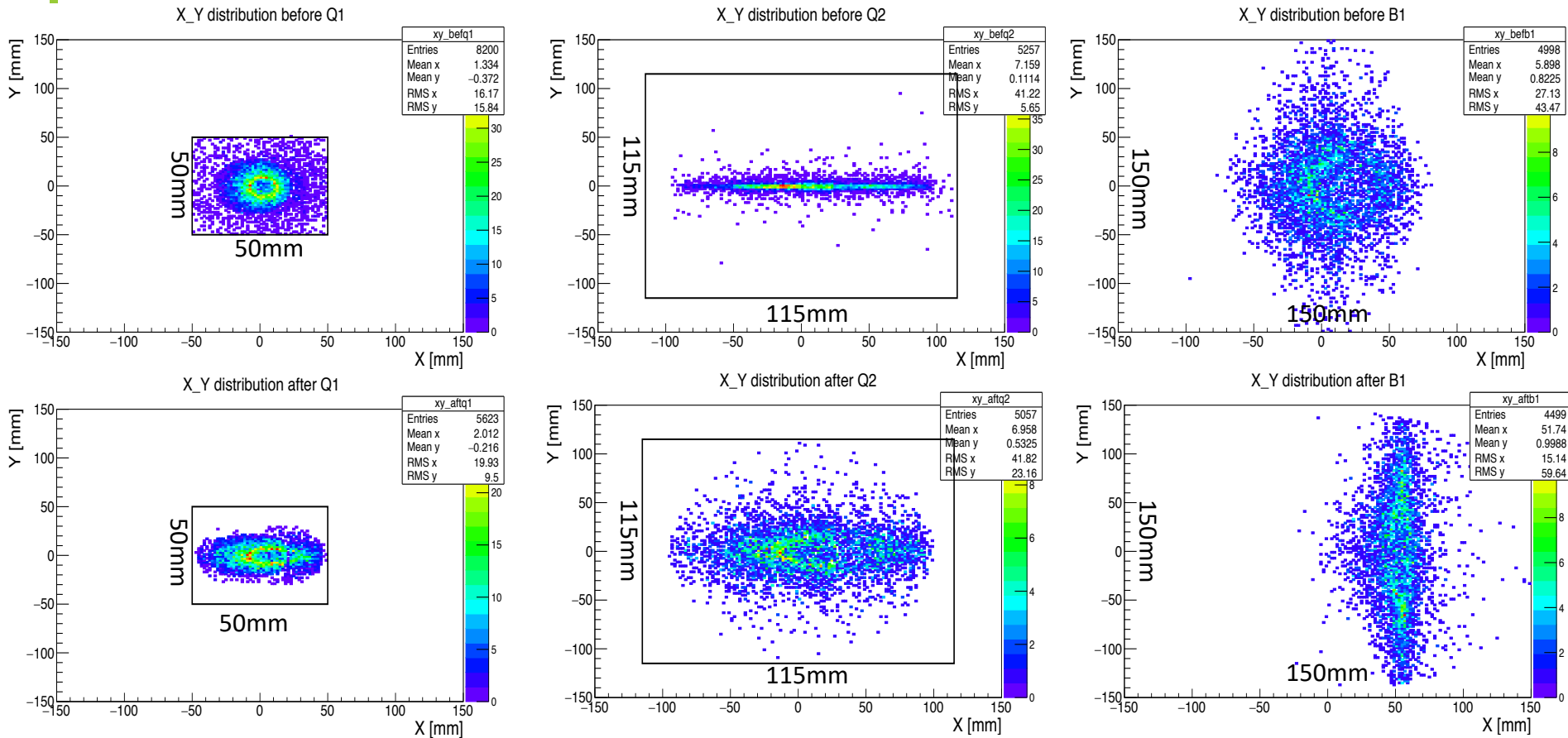
Acceptance



	Center_ z [m]	Center_ x [m]	R_{in} [m]	Length [m]
B0PF	5.5999	0.1237	0.200	1.200
Q1PF	7.5498	0.1695	0.042	1.500
Q2PF	12.1999	0.2803	0.105	2.400
B1PF	17.9002	0.4226	0.135	3.000

- ❖ 20% loss of acceptance in ~ 0 theta from aft q1 to bef q2.
- ❖ Roman Pots have much less acceptance.

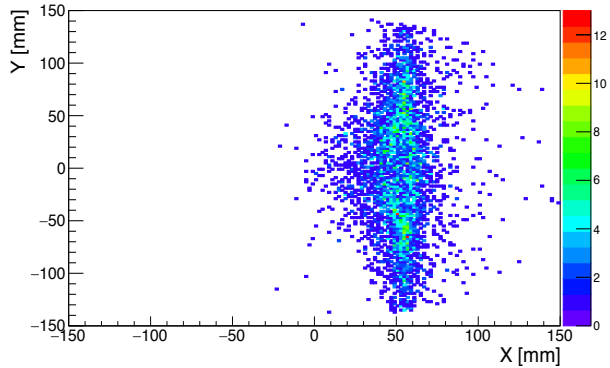
X_Y distribution



There isn't have 10σ away cut here, because we can not put a real detector in the main beam

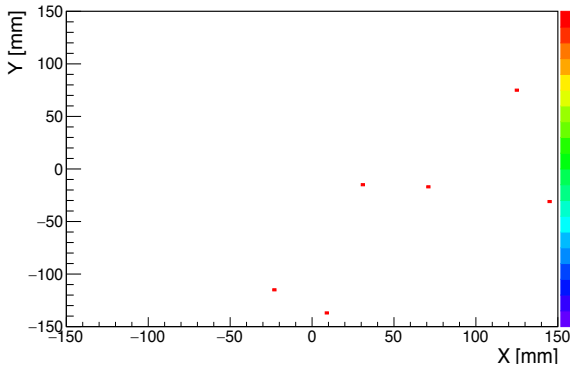
X-Y distribution after B1

X_Y distribution after B1

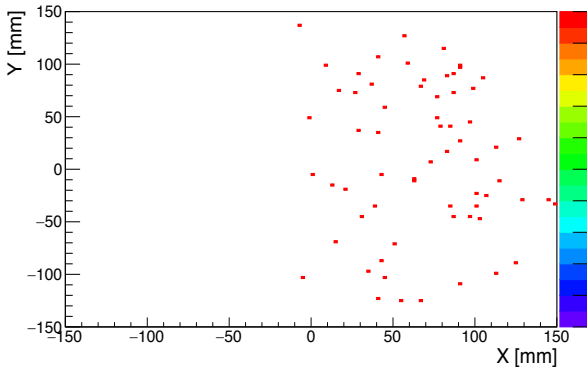


As the p decreases, the rigidity goes down,
and the proton bends more

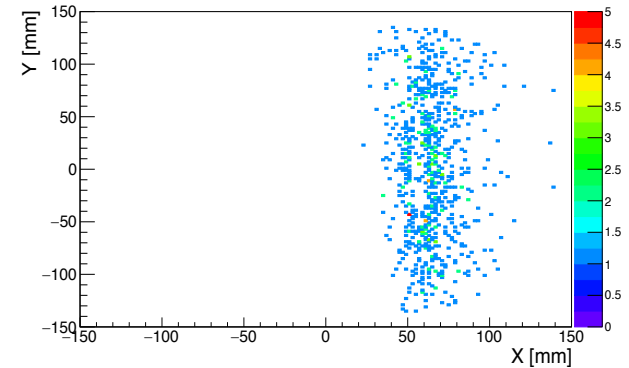
X_Y distribution after B1 for momentum of 0_60 GeV



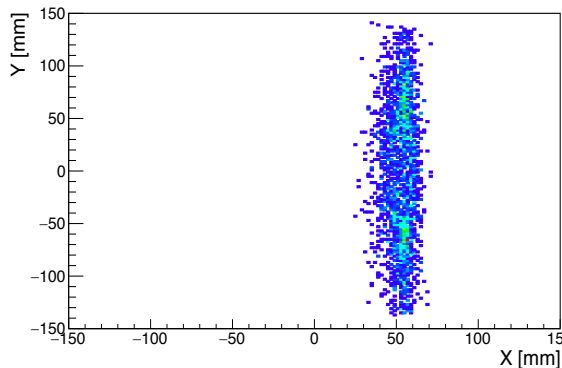
X_Y distribution after B1 for momentum of 60_80 GeV



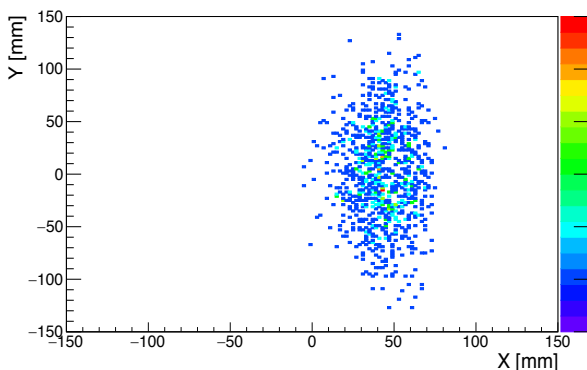
X_Y distribution after B1 for momentum of 80_100 GeV



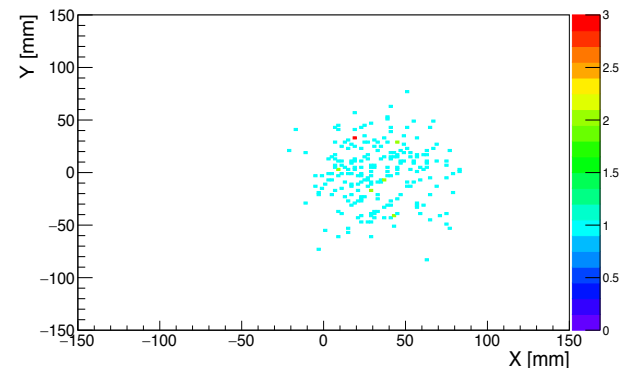
X_Y distribution after B1 for momentum of 100_120 GeV



X_Y distribution after B1 for momentum of 120_140 GeV

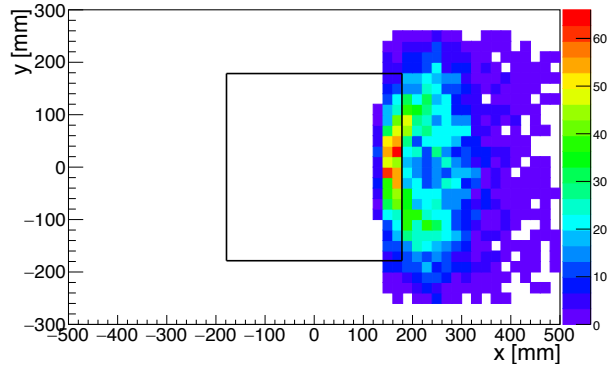


X_Y distribution after B1 for momentum of 140_200 GeV



X-Y distribution at Roman Pots position

X_Y distribution at Roman Pots

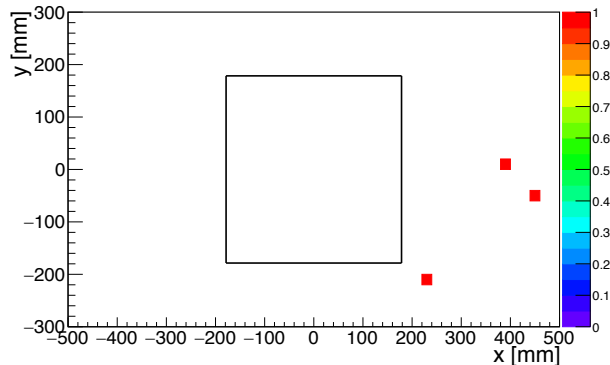


Roman Pots: 28.1m

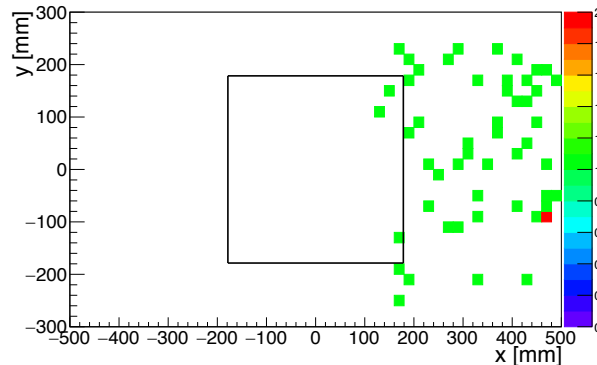
Possible real size: 375mm * 375mm

Roman pots have not enough acceptance, need a different detector solution

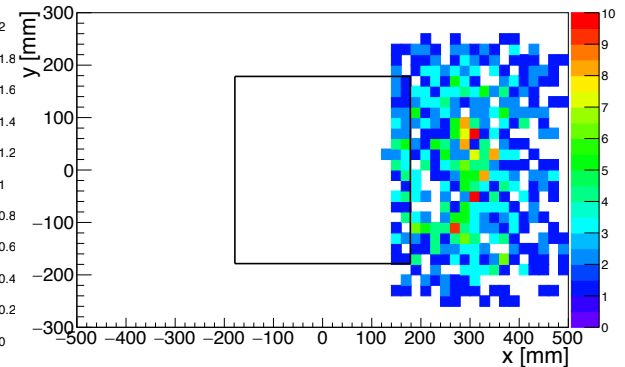
X_Y distribution for momentum of 0_60 GeV at Roman Pots



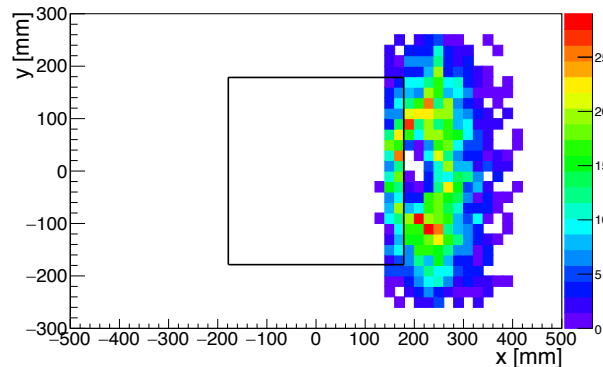
X_Y distribution for momentum of 60_80 GeV at Roman Pots



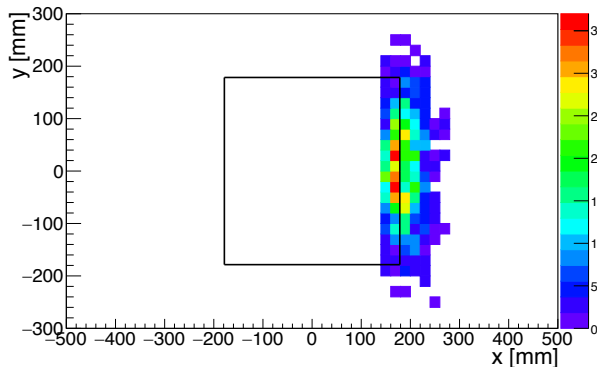
X_Y distribution for momentum of 80_100 GeV at Roman Pots



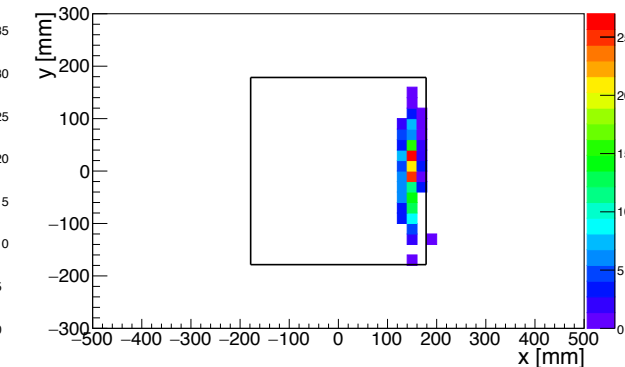
X_Y distribution for momentum of 100_120 GeV at Roman Pots



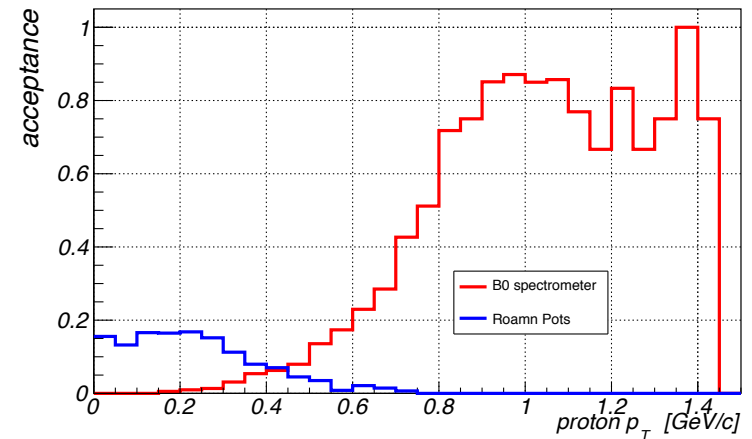
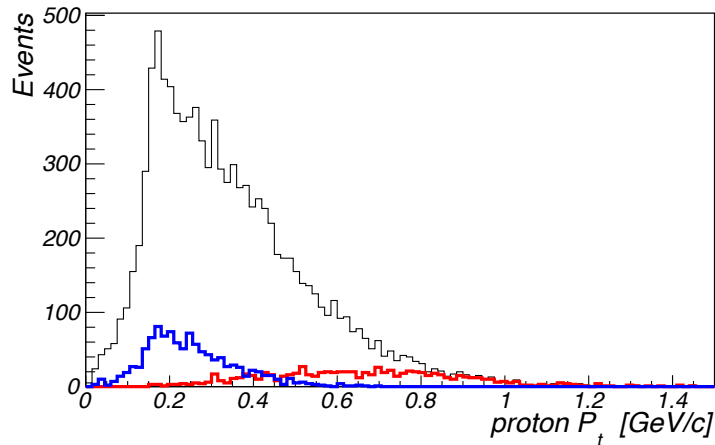
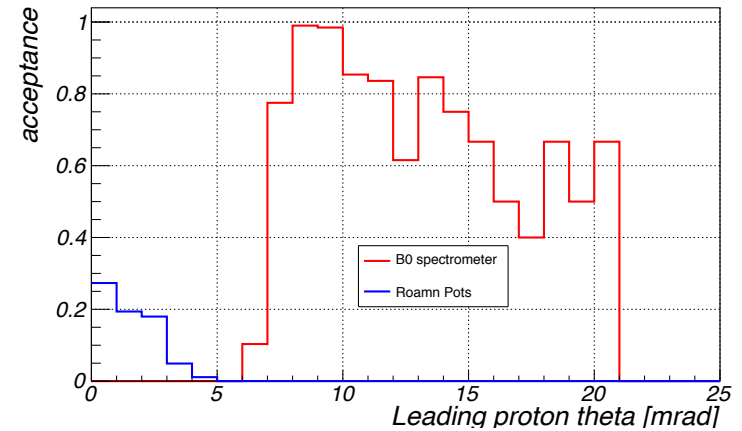
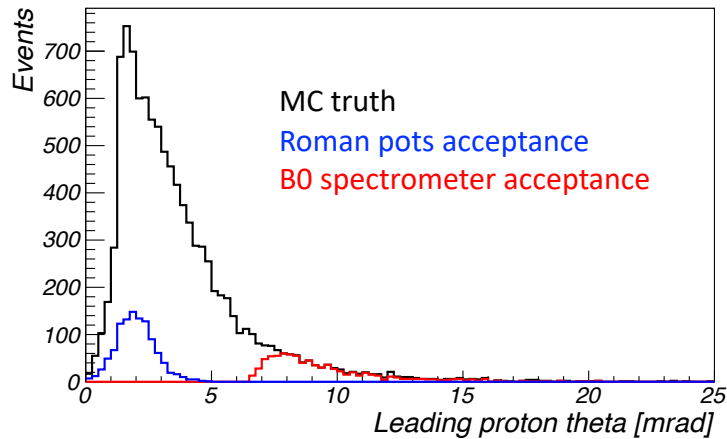
X_Y distribution for momentum of 120_140 GeV at Roman Pots



X_Y distribution for momentum of 140_200 GeV at Roman Pots



Proton acceptance in ePb



MC generated: 10K; Roman Pots accepted: 1.1K

It's very little statistics at B0, the Roman Pots acceptance for proton is poor

Summary

- ❑ We studied the acceptance of protons in B0 and Roman Pots in ep and ePb collisions.
- ❑ The Roman Pots acceptance for protons from Pb-decay is poor, because of the different rigidity between proton and Pb.
- ❑ For ePb, we need to have an additional detector, which recovers this loss.

The comparison between BeAGLE and E665

WAN CHANG

Events information

Data sample:

$\mu^+ + \text{Xe}$

Beam momentum:

490 GeV \times 0 GeV

$0.1 < y < 0.85$

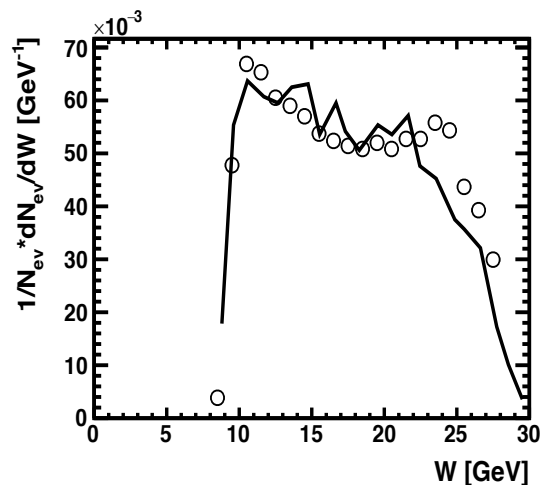
$1.0 < Q^2 < 100$

$0.0035 \text{ rad} < \theta < 6.29 \text{ rad}$

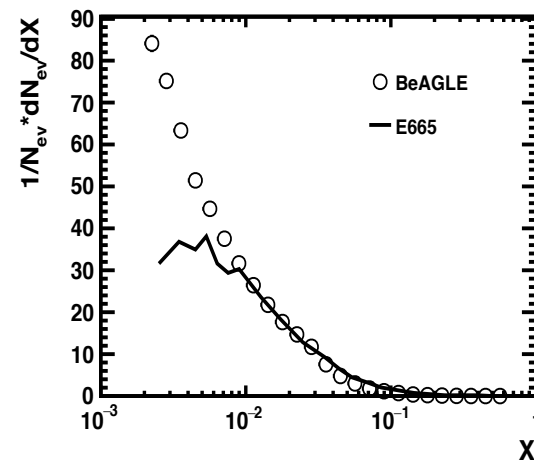
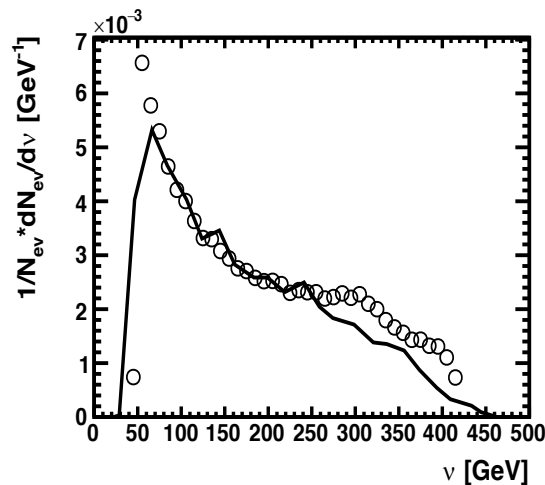
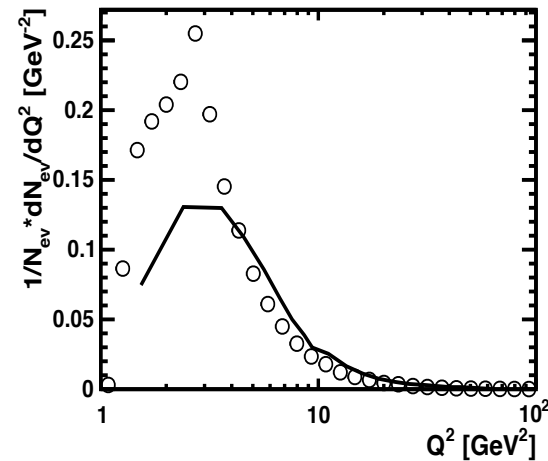
$8 < W < 30 \text{ GeV}$

$X > 0.002$

40K events

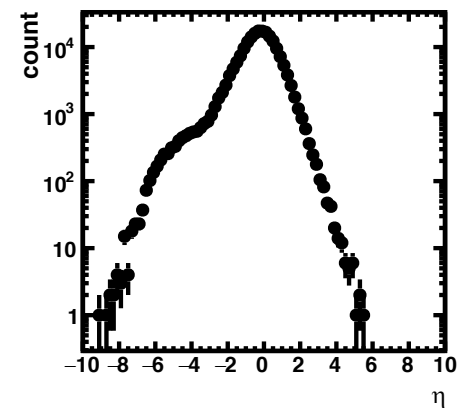
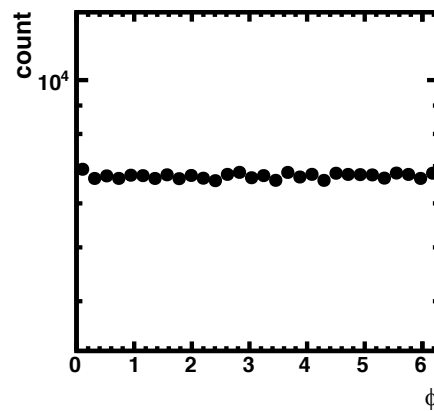
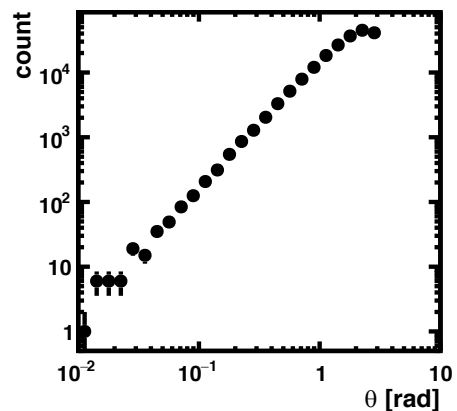
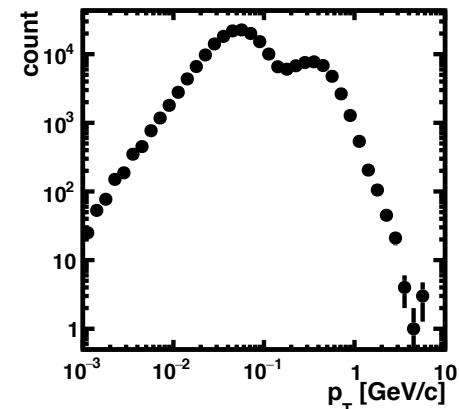
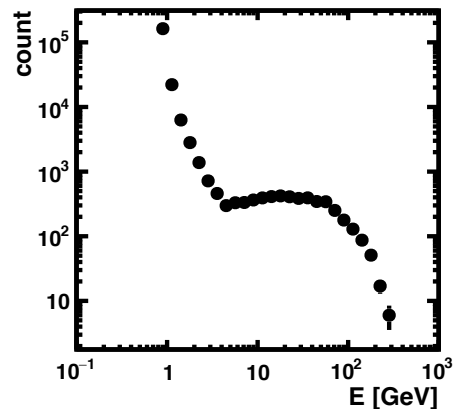
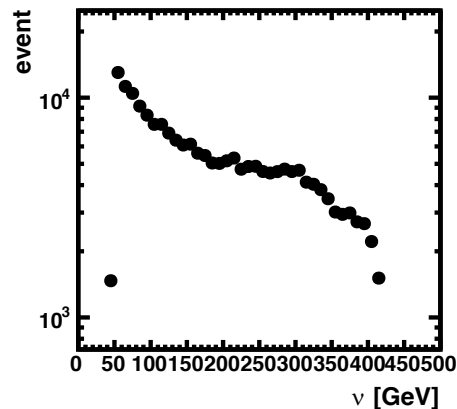


Z. Phys. C 61, 179-198(1994)



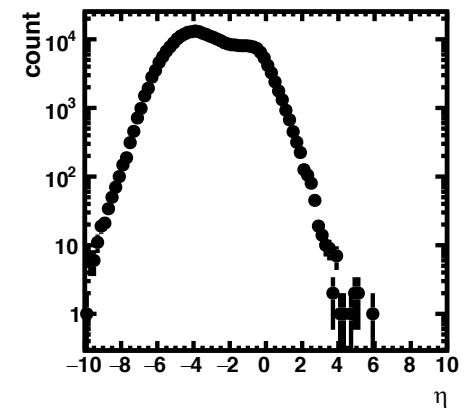
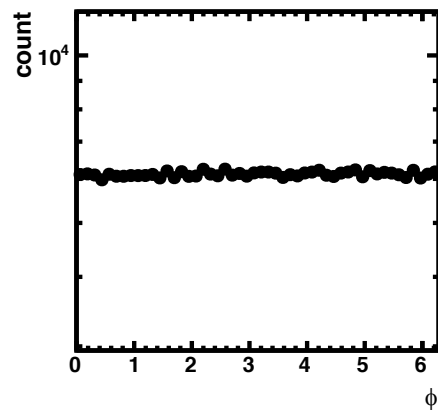
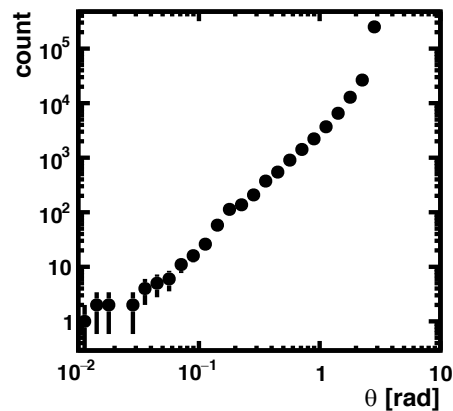
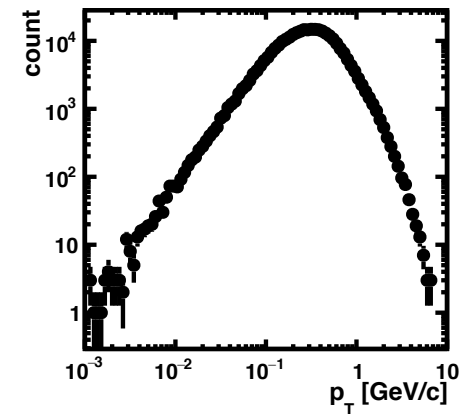
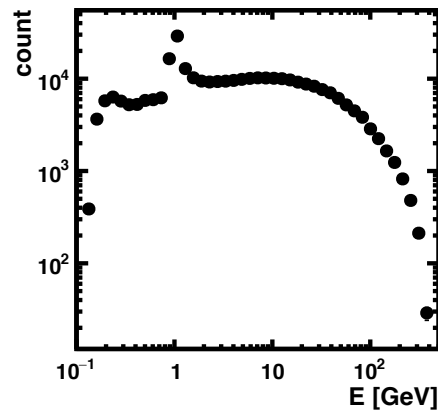
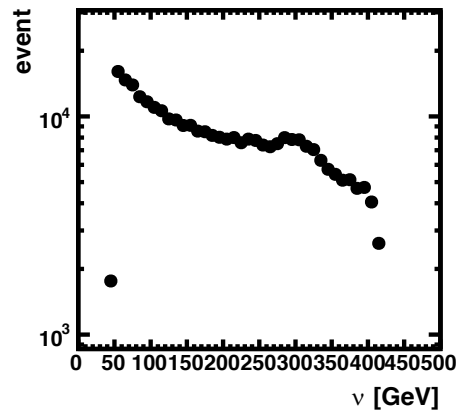
Neutron

Raw distribution of $\nu, E, p_T, \theta, \phi, \eta$ for **neutrons** in the **lab** frame:



Hadron

Raw distribution of $\nu, E, p_T, \theta, \phi, \eta$ for **hadrons** in the **lab** frame:



To do ...

To do the comparison on:

- Event kinematic (X, y, Q^2, v, W). VS Z. Phys. C 61, 179-198(1994)
- Track distribution($v, E, p_T, \theta, \phi, \eta, \frac{1}{N_{ev}} \frac{dN}{dy^*}, \dots$) VS Z. Phys. C 61, 179-198(1994)
- Multiplicity distributions($P(n), \bar{n}, \text{VS } W^2$) VS Z. Phys. C 61, 179-198(1994)
- The average total hadronic net charge $\langle Q_T \rangle$, the average number of grey tracks $\langle n_g \rangle$, and the difference of average charged backward multiplicities $\langle n_B \rangle_{\mu Xe} - \langle n_B \rangle_{\mu D}$
VS Z. Phys. C 65, 225-244(1995)
- Multiplicity ratio $(R_A^h = \frac{\left(\frac{N_h(v, Q^2, p_T^2, z)}{N_l(v, Q^2)} \right)_A}{\left(\frac{N_h(v, Q^2, p_T^2, z)}{N_l(v, Q^2)} \right)_D})$ VS HERMES collaboration
- And so on.....

E665 collaboration:

Z. Phys. C 61, 179-198(1994)

Z. Phys. C 65, 225-244(1995)

Phys. Rev. Lett. 74, 5198

HERMES collaboration:

arXiv:0906.2478

arXiv:hep-ex/0510030

arXiv:0704.3270

arXiv:1107.3496